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# Electronic structure of the d<sup>1</sup> bent-metallocene Cp<sub>2</sub>VCl<sub>2</sub>: A photoelectron and density functional study

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In tribute to F. Albert Cotton for not just his teachings, but also for the excitement of discovery and new understanding in chemistry he brought to us and many others.

#### Abstract

The  $Cp_2VCl_2$  molecule is a prototype for bent-metallocene complexes with a single electron in the metal d shell, but experimental measure of the binding energy of the d electron by photoelectron spectroscopy eluded early attempts due to apparent decomposition in the spectrometer to  $Cp_2VCl$ . With improved instrumentation, the amount of decomposition is reduced and subtraction of ionization intensity due to  $Cp_2VCl$  from the  $Cp_2VCl_2/Cp_2VCl$  mixed spectrum yields the  $Cp_2VCl_2$  spectrum exclusively. The measured ionization energies provide well-defined benchmarks for electronic structure calculations. Density functional calculations support the spectral interpretations and agree well with the ionization energy of the d<sup>1</sup> electron and the energies of the higher positive ion states of  $Cp_2VCl_2$ . The calculations also account well for the trends to the other Group V bent-metallocene dichlorides  $Cp_2Ncl_2$  and  $Cp_2TaCl_2$ . The first ionization energy of  $Cp_2VCl_2$  is considerably greater than the first ionization energies of the second- and third-row transition metal analogues.

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# 1. Introduction

The chemistry of bent-metallocene complexes has been a major factor in the development of organometallic chemistry over the last 50 years [1-5], finding applications in diverse fields ranging from polymer catalysis [4,6,7] to cancer research [8-11]. The electron configuration and subsequent reactivity of the metal center can be fine-tuned with the choice of metal, cyclopentadienyl substituents, and additional coordinating ligands. Elucidating the electronic structure of bent-metallocenes is important in understanding the chemical behavior and physical properties of these complexes. Photoelectron spectroscopy provides an experimental measure of the orbital energetics and interac-

\* Corresponding author. E-mail address: dlichten@email.arizona.edu (D.L. Lichtenberger). tions and at this point there is a general understanding of the photoelectron spectra of simple metallocenes [7,12,13]. One remaining point of contention has been the reported photoelectron spectrum of  $Cp_2VCl_2$ , which has been an outlier from all other reported metallocene photoelectron results. This has been unfortunate because  $Cp_2VCl_2$  is a prototype for d<sup>1</sup> bent-metallocene complexes, having a first-row transition metal, relatively simple ligands, and sufficient symmetry to be a foundational model for other studies.

The electronic structure of Cp<sub>2</sub>VCl<sub>2</sub> has previously been investigated by electron paramagnetic resonance (EPR) [14–17] and photoelectron spectroscopy (PES) [15]. EPR experiments concluded that the unpaired metal electron resides in an orbital composed primarily of  $d_{z^2}$  character in the VCl<sub>2</sub> plane as depicted in the coordinates below, with some admixture of  $d_{x^2-y^2}$  character. Photoelectron studies of Cp<sub>2</sub>VCl<sub>2</sub> and (C<sub>5</sub>H<sub>4</sub>Me)<sub>2</sub>VCl<sub>2</sub> [15,18] revealed spectra

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that were similar in structure to those of their Ti analogues, but with two low energy ionization features. Only one low energy ionization was expected for the single d electron in the vanadium complexes. The possibility of decomposition was investigated and considered less likely for the (C5H4Me)2VCl2 molecule. The lowest energy feature of the spectrum obtained from the (C<sub>5</sub>H<sub>4</sub>Me)<sub>2</sub>VCl<sub>2</sub> sample at 6.60 eV was tentatively assigned to the metal d<sup>1</sup> ionization. The second feature at 7.20 eV was hypothesized to be a triplet component of the unpaired  $d^1$  electron exchange splitting with the Cl  $p\pi$  orbitals. The singlet component of this ionization could then contribute to broadening of the next band, rather than appear as a separate band. This assignment was called into doubt when it was found that the related Nb and Ta complexes had only one low energy ionization consistent with their d<sup>1</sup> configurations [12,18]. Subsequent photoelectron studies of Cp<sub>2</sub>VCl [19], which has a high-spin  $d^2$  configuration, showed bands at 6.80 and 7.40 eV. The spectrum of this monochloride complex appeared similar to the reported spectrum of the dichloride complex, suggesting that Cp<sub>2</sub>VCl<sub>2</sub> may be decomposing to Cp<sub>2</sub>VCl during the photoelectron experiment. Cp<sub>2</sub>VCl<sub>2</sub> is known to undergo thermal- [15] and photoelectrochemical decomposition [20] to Cp<sub>2</sub>VCl during other experimental conditions.



Recently, we investigated the electronic structure of bent vanadocene dithiolates as minimum molecular models of the d<sup>1</sup> electron configuration of sulfite oxidizing enzymes using photoelectron spectroscopy and density functional theory [21]. This study prompted us to reinvestigate the electronic structure of the simplest d<sup>1</sup> bent-metallocene, Cp<sub>2</sub>VCl<sub>2</sub>, in more detail to aid in the interpretation of the vanadocene dithiolate spectra. Improved gas-phase photoelectron spectroscopy techniques, comparison to the photoelectron spectra of Cp<sub>2</sub>VCl and other Group V metallocene dichlorides, along with DFT calculations, have allowed us to isolate and reassign the photoelectron spectrum of Cp<sub>2</sub>VCl<sub>2</sub> for determination of the ionization energy of the d<sup>1</sup> electron.

### 2. Results and discussion

# 2.1. Photoelectron spectroscopy

The low energy valence regions of the gas-phase photoelectron (PES) spectra collected from samples of  $Cp_2VCl$  and Cp<sub>2</sub>VCl<sub>2</sub> are presented in Fig. 1. The spectra obtained using typical data collection techniques are in agreement with spectra previously reported in the literature [15,19]. There is a large similarity between the two spectra obtained from the Cp<sub>2</sub>VCl and Cp<sub>2</sub>VCl<sub>2</sub> samples, Fig. 1A and B respectively, which would not be expected given the anticipated relative ionizations from the high-spin d<sup>2</sup> metal center of  $Cp_2VCl$  versus the d<sup>1</sup> metal center of  $Cp_2VCl_2$ . When a spectrum of the Cp<sub>2</sub>VCl<sub>2</sub> sample is recorded using sample entry procedures designed to minimize decomposition (see Section 4), the overall shape of the spectrum changes dramatically (Fig. 1C). Most importantly, the relative intensity of the ionization at 6.8 eV decreases substantially relative to the ionization at 7.4 eV. In addition, ionization intensity in the region labeled L3 in the spectrum of Cp<sub>2</sub>VCl decreases while the ionization intensity in the region of L2 increases. The similarities in the spectra 1A and 1B and the differences between these and spectrum 1C are strong evidence that Cp<sub>2</sub>VCl<sub>2</sub> is decomposing to Cp<sub>2</sub>VCl in the spectrometer. The decomposition has been minimized, but not completely eliminated, as evidenced by the small amount of ionization intensity remaining at 6.8 eV in Fig. 1C. The remaining ionization intensity in spectrum 1C from Cp<sub>2</sub>VCl cannot be due to residual Cp<sub>2</sub>VCl in the Cp<sub>2</sub>VCl<sub>2</sub> sample because the sublimation temperature of Cp<sub>2</sub>VCl is approximately 100 °C lower than that of Cp<sub>2</sub>VCl<sub>2</sub> (see Section 4) and would be observed first as the temperature of the sample cell rises. Also, this contamination cannot be due to photoelectrochemical dehalogenation [20] of the solid sample since the sample is in the



Fig. 1. Low energy He I spectra showing the extraction of the  $Cp_2VCl_2$  spectrum from the data: (A) spectrum of  $Cp_2VCl_2$  (B) spectrum of  $Cp_2VCl_2$  sample loaded directly into ionization chamber, (C) spectrum of  $Cp_2VCl_2$  sample effused directly into photon beam from a quartz crucible, (D) subtraction of A from C to give difference  $Cp_2VCl_2$  spectrum.

direct path of the ionization source during the experiment and Cp<sub>2</sub>VCl is not seen at lower temperatures, nor were ionizations found corresponding to the photochemical product Cl<sub>2</sub>. Therefore, it is hypothesized that use of the quartz crucible and the short path to the photon beam reduces surface interactions of Cp<sub>2</sub>VCl<sub>2</sub> with the aluminum sample cell, which may lead to reduction and dehalogenation of Cp<sub>2</sub>VCl<sub>2</sub>. It was noted that, during the experiment, a sharp doublet-ionization occurs at 12.75 and 12.82 eV, coincident with the first ionizations of HCl. The ionization intensity due to Cp<sub>2</sub>VCl in Fig. 1C is removed by subtracting the Fig. 1A spectrum, appropriately scaled to the ionization of Cp<sub>2</sub>VCl at 6.8 eV in Fig. 1C, to yield the spectrum shown in Fig. 1D. The ionizations shown in Fig. 1D are taken to be those of Cp<sub>2</sub>VCl<sub>2</sub>. The subtraction of the Cp<sub>2</sub>VCl spectrum leads to a shape (width and symmetry) for the first ionization of Cp<sub>2</sub>VCl<sub>2</sub> at 7.40 eV that is indicative of a single ionization [22].

Fig. 2 shows that the resultant  $Cp_2VCl_2$  spectrum bears strong similarities to the He I photoelectron spectra of  $Cp_2NbCl_2$  and  $Cp_2TaCl_2$  [12]. The ionization region of ~8–9.8 eV is particularly diagnostic, because this region contains ionizations that are symmetry combinations of the  $p\pi$  orbitals of the two Cl atoms with the metal. The similar ionization pattern in this region indicates similar  $MCl_2$  structure for the three molecules. The main difference between the spectra of  $Cp_2MCl_2$  (M = V, Nb, Ta) (Fig. 2) is the energy of the metal d<sup>1</sup> ionization, which can be ascribed to reduction of the effective nuclear charge on going from 3d to 4d to 5d, reflecting the shielding of the unpaired electron by the different atoms.

Table 1 shows the analytical representation of the ionizations for  $Cp_2VCl_2$  and  $Cp_2VCl$ . This table includes the



Fig. 2. He I photoelectron spectra of Cp<sub>2</sub>VCl<sub>2</sub>, Cp<sub>2</sub>NbCl<sub>2</sub> and Cp<sub>2</sub>TaCl<sub>2</sub>.

change in band area with change in ionization source from He I to He II for Cp<sub>2</sub>VCl. Green et al. have previously assigned the photoelectron spectrum of Cp<sub>2</sub>VCl based on a He I/He II comparison and their assignments are reflected in Table 2 [19]. The predominant Cl  $p\pi$  character mentioned above in ionizations labeled L1 and L2 is evidenced by their decrease in ionization intensity relative to the metal-based ionizations from He I to He II excitation.

## 2.2. Computational results

A central question concerns the ability of electronic structure calculations to account for the geometric structures and the electron energies and distributions of openshell bent-metallocene complexes of the early transition metals. The calculated structures of Cp<sub>2</sub>VCl<sub>2</sub> and Cp<sub>2</sub>VCl are compared with the reported structures determined from X-ray crystallography in Table 3 [23,24]. The structures agree reasonably well with the largest deviation being the V–Cl bond distance in Cp<sub>2</sub>VCl<sub>2</sub> which is underestimated by about 0.04 Å. The ground states of these molecules are calculated as an unrestricted doublet and triplet, respectively for Cp<sub>2</sub>VCl<sub>2</sub> and Cp<sub>2</sub>VCl. Spin contamination is minimal, giving a value of 0.78 compared to the ideal value for *s*(*s* + 1) of 0.75 for Cp<sub>2</sub>VCl<sub>2</sub> and 2.05 compared to the ideal value of 2.00 for Cp<sub>2</sub>VCl.

DFT calculations were also used to model the electron density in the highest occupied molecular orbital (HOMO) for comparison with EPR studies. The trend of the Cl–M– Cl angle to decrease from Ti to V to Mo for the d<sup>0</sup> to d<sup>1</sup> to d<sup>2</sup> configurations and the single crystal EPR studies on the d<sup>1</sup> complexes led to the conclusion that the HOMO of the V and Mo compounds is primarily along an axis normal to the plane bisecting the Cl–M–Cl angle. Lowering the symmetry of Cp<sub>2</sub>VCl<sub>2</sub> from  $C_{2v}$  to  $C_s$  allows this axis to be labeled z, as illustrated in Section 1. If the ligand contributions

Table 1 Analytical representation of the ionization features of  $Cp_2VCl_2$  and  $Cp_2VCl$  obtained with He I and He II photon sources

Band	IE <sup>a</sup> (eV)	Band width <sup>b</sup>	Relativ	Relative area <sup>c</sup>		
		High (eV)	Low (eV)	He I	He II	
$Cp_2VCl_2$	,					
M1	7.40	0.30	0.30	1		
L1	8.47	0.47	0.38	5.42		
L2	8.85	0.55	0.26	8.80		
$Cp_2VCl$						
M1	6.78	0.33	0.23	1	1	
M2	7.40	0.22	0.12	0.78	0.97	
L1	8.33	0.39	0.39	3.67	1.41	
L2	9.04	0.36	0.36	2.32	1.41	
L3	9.45	0.54	0.49	5.92	4.22	

<sup>a</sup> Vertical ionization energy defined as the position of the asymmetric Gaussian peak modeling the band.

<sup>b</sup> Widths of the asymmetric Gaussian peak modeling the high and low ionization energy sides of the band.

Band areas are relative to an area of 1 for M1.

Table 2	
Experimental and calculated ionization energies (eV) for Cp <sub>2</sub> VCl <sub>2</sub> and Cp <sub>2</sub> VCl	

Band	$IE^{a}$	Assignment	Kohn-Sham		$\Delta SCF^{c}$	
			Orbital energy <sup>b</sup>			
$Cp_2VCl_2$						
-			α-Spin	β-Spin	Singlet	Triplet
M1	7.40	$V d_{z^2}$	-4.83(7.40)		7.27	_
L1	8.47	Cl pπ	-5.81(8.38)	-5.78(8.35)	8.20	8.15
L2	8.85	Cl pπ	-5.96 (8.53)	-5.88 (8.45)	8.49	8.44
$Cp_2VCl$						
-			α-Spin	β-Spin	Doublet	Quartet
M1	6.78	$V d_{xz}$	-4.24 (6.78)		6.85	
M2	7.40	$V d_{z^2}$	-4.87 (7.41)		7.56	
L1	8.33	Cl pπ	-5.64(8.18)	-5.59 (8.13)	8.33	8.18
L2	9.04	Cl pπ	-6.65 (9.19)	-6.34 (8.88)	9.27	8.91
L3	9.45	$Cp p\pi + Cl p\pi$	-7.17 (9.71)	-6.93 (9.47)	9.50	9.24

<sup>a</sup> Experimental vertical ionization energy in eV.

<sup>b</sup> The values in parentheses are the orbital energies shifted by an amount such that the first orbital energy matches the first ionization energy.

<sup>c</sup> The difference in total energy between the molecular ground state and the molecule with an electron removed from the specified orbital.

Table 3							
Comparison	of	the	calculated	geometries	with	the	crystallographic
molecular str	ucti	ires f	for Cp <sub>2</sub> VCl <sub>2</sub>	and Cp <sub>2</sub> VC	la		

	Calculated	Experimental
$Cp_2VCl_2^b$		
Cp <sub>C-C</sub>	$1.403 - 1.428 (1.415)^{c}$	1.371-1.465 (1.418)
		1.394-1.457 (1.424)
M-Cp <sub>centroid</sub>	1.952, 1.948	1.983, 1.967
		1.971, 1.986
M–C	2.231-2.306 (2.269)	2.281-2.368 (2.314)
		2.284-2.342 (2.315)
M–Cl	2.371	2.409, 2.418
		2.410, 2.411
Cl-M-Cl	88.74°	86.6°
		87.1°
$Cp_2VCl^d$		
Cp <sub>C-C</sub>	1.406-1.420 (1.413)	1.379-1.408 (1.394)
M-Cp <sub>centroid</sub>	1.901, 1.901	1.946, 1.944
M–C	2.221-2.277 (2.250)	2.261-2.296 (2.278)
M-Cl	2.367	2.390

<sup>a</sup> All distances are given in Å.

<sup>b</sup> Crystallographic data taken from Tzavellas et al. [23].

<sup>c</sup> Average distance.

<sup>d</sup> Crystallographic data taken from Fieselmann et al. [24].

are neglected then  $|\Psi_{\text{HOMO}}\rangle = a|\mathbf{d}_{z^2}\rangle + b|\mathbf{d}_{x^2-y^2}\rangle$ . Previous Fenske-Hall [25] molecular orbital percent characters in the singly-occupied orbital were calculated to be 20.5/1 for  $\mathbf{d}_{z^2}/\mathbf{d}_{x^2-y^2}$  in Cp<sub>2</sub>VCl<sub>2</sub>, which agrees well with the ratio obtained from the EPR data for (C<sub>6</sub>H<sub>5</sub>Me)<sub>2</sub>VCl<sub>2</sub> of 20.0/1 [15]. Using current DFT methods, the singly-occupied orbital of Cp<sub>2</sub>VCl<sub>2</sub> is calculated to be 60.8%  $\mathbf{d}_{z^2}$  and 3.44%  $\mathbf{d}_{x^2-y^2}$  giving a ratio of 17.7:1. This ratio is similar to the orbital percent characters calculated using the Fenske–Hall method and the observed orbital parameters calculated from the EPR experiment.

The energies of the ionizations observed for  $Cp_2VCl_2$ and  $Cp_2VCl$  by photoelectron spectroscopy are compared with those calculated by the  $\Delta$ SCF method and with the Kohn-Sham orbital energies (Table 2). The calculated Kohn-Sham orbital energies from the DFT calculations are related to the photoelectron ionization energies by applying a corollary to Koopmans' theorem from Hartree-Fock calculations [26]. Koopmans' theorem has intrinsic uncertainty including electron relaxation, electron correlation, relativistic effects, vibronic coupling, and differences in spin states [27-34], which can lead to poor correlation between the calculated orbital energies and the ionization potentials. A central principle of the Kohn-Sham orbital model is that the negative of the Kohn-Sham orbital energy for the HOMO, calculated with the correct functional and basis, will exactly match the first ionization energy of the molecule, but in current practice these energies differ substantially. Nonetheless, the relative energies of the Kohn-Sham orbitals can be useful in interpreting the pattern of ionizations.  $\Delta$ SCF calculations were performed to account for the deficiencies in the frozen orbital approximation and to get information on the coupling of electron spins.

The ground state of neutral Cp<sub>2</sub>VCl<sub>2</sub> is  $s = \frac{1}{2}$  (V<sup>4+</sup>, d<sup>1</sup>, <sup>2</sup>A') while the ground state of Cp<sub>2</sub>VCl is s = 1 (high-spin  $V^{3+}$ ,  $d^2$ ,  ${}^{3}A''$ ). Fig. 3 shows the contour plots for the spin  $\alpha$  and  $\beta$  orbitals of Cp<sub>2</sub>VCl<sub>2</sub> that correspond to the ionizations evaluated in this study by PES. For Cp<sub>2</sub>VCl<sub>2</sub>, ionization from the singly-occupied HOMO will lead to a singlet state, while ionization from the doubly-occupied ligandbased orbitals leads to both singlet and triplet state configurations with coupling to the d<sup>1</sup> electron. For Cp<sub>2</sub>VCl, ionization from the singly-occupied HOMO and HOMO-1 lead to doublet states, and ionization from the doubly-occupied ligand-based orbitals will lead to either doublet or quartet state configurations. Table 2 shows the calculated values for the singlet/triplet states of Cp<sub>2</sub>VCl<sub>2</sub> and doublet/quartet states of Cp<sub>2</sub>VCl using the  $\Delta$ SCF method. For Cp<sub>2</sub>VCl<sub>2</sub>, the calculated singlet/triplet state separation is only 0.05 eV for both ionizations L1



Fig. 3. Spin correlation diagram for  $Cp_2VCl_2$  showing  $\alpha$ - and  $\beta$ -spin molecular orbital character and electron occupations.

and L2, and therefore ionizations to these individual states are not likely to be resolved in the photoelectron spectrum. For Cp<sub>2</sub>VCl the calculated doublet/quartet state separation for band L1 is 0.15 eV, still less than the 0.4 eV width of the unresolved vibrational broadening of the band. For bands L2 and L3, the calculated doublet/quartet state separations are 0.36 and 0.26 eV, respectively; however, these bands are now in a region of overlapping ionizations that preclude observation of individual states.

 $\Delta$ SCF calculations for Cp<sub>2</sub>VCl<sub>2</sub> predict the singly-occupied HOMO ionization to be at 7.27 eV, consistent with the first ionization of Cp<sub>2</sub>VCl<sub>2</sub> at 7.40 eV (Table 2). The calculated  $\Delta$ SCF ionizations for L1 and L2 are also similar with the HOMO-1 and HOMO-2 ionizations, and follow the ionization trend. The shifted Kohn-Sham orbital energies for L1 and L2 are also similar to the observed ionization energies.  $\Delta$ SCF calculations were also performed on Cp<sub>2</sub>VCl to confirm that removing an electron from the triplet ground state d orbital configuration will give a metal ionization that coincides with removing an electron from the HOMO of  $Cp_2VCl_2$ . Table 2 shows that removing an electron from the HOMO-1 of  $Cp_2VCl$  gives a  $\Delta SCF$  calculated ionization energy of 7.56 and a shifted Kohn-Sham orbital energy of 7.41, coincident with the HOMO ionization of  $Cp_2VCl_2$ . Both the  $\Delta SCF$  method and shifted Kohn-Sham orbital energies follow similar trends to the observed ionization energies for Cp<sub>2</sub>VCl. These calculations support that subtraction of the Cp<sub>2</sub>VCl spectral contribution from the mixed Cp<sub>2</sub>VCl<sub>2</sub>/Cp<sub>2</sub>VCl spectrum is effective for revealing the  $Cp_2VCl_2$  spectrum. The first ionization of  $Cp_2NbCl_2$  and  $Cp_2TaCl_2$  was also calculated by the  $\Delta$ SCF method to be 6.65 and 6.37 eV, almost exactly coincident to the observed ionization energies of 6.69 and 6.39 eV [12].

# 3. Conclusions

We have revisited the gas-phase photoelectron spectrum of Cp<sub>2</sub>VCl<sub>2</sub> and used improved collection techniques to reduce the Cp<sub>2</sub>VCl decomposition product from the Cp<sub>2</sub>VCl<sub>2</sub> He I spectrum. Subtraction of ionization intensity due to the Cp<sub>2</sub>VCl decomposition product leads to a  $Cp_2VCl_2$  spectrum that is similar in appearance to other Group V bent-metallocene dichlorides. Density functional theory calculations support the assignments made for the Cp<sub>2</sub>VCl<sub>2</sub> spectrum and comparison to the calculated ionization energies for Cp<sub>2</sub>VCl also confirm that the d<sup>1</sup> ionization of Cp<sub>2</sub>VCl<sub>2</sub> coincides with one of the metal ionizations of Cp<sub>2</sub>VCl, further supporting the superposition of the Cp<sub>2</sub>VCl and Cp<sub>2</sub>VCl<sub>2</sub> spectra. The open-shell calculations on these Group V bent-metallocenes show good reliability in calculating the geometric structures and electron distributions and give good account of the ionization energies using either the shifted Kohn-Sham orbital energies or the  $\Delta$ SCF method.

#### 4. Experimental methods

# 4.1. Photoelectron spectroscopy

Samples of Cp<sub>2</sub>VCl (97%) and Cp<sub>2</sub>VCl<sub>2</sub> (95%) were purchased through Aldrich and used as received. Photoelectron spectra were recorded using an in-house instrument built around a 36 cm hemispherical analyzer (McPherson), custom-designed sample cells and detection and control electronics. The electron detection and instrument operation are interfaced to a National Instruments PCIe-6259 multi-function data acquisition card and custom software. Argon was used as an internal calibrant during data collection and the instrument resolution (measured using FWHM of the argon  ${}^{2}P_{3/2}$  peak) was 0.020-0.030 eV. The sublimation temperatures (at  $10^{-5}$  Torr, monitored using a "K" type thermocouple passed through a vacuum feed-through and attached directly to the sample cell) were 195-205° for Cp<sub>2</sub>VCl<sub>2</sub> and 90-120° for Cp<sub>2</sub>VCl. For Cp<sub>2</sub>VCl<sub>2</sub>, a crushed crystalline sample was placed in a quarter-inch diameter quartz crucible that was inserted into the ionization chamber of the sample cell directly below the photon beam [35,36].

In the photoelectron spectra, the vertical length of each data mark represents the experimental variance of that point. The valence ionization bands are represented analytically with the best fit of asymmetric Gaussian peaks [22]. The number of peaks used in a fit was based solely on the features of a given band profile. The peak positions are reproducible to about  $\pm 0.02 \text{ eV} (\approx 3\sigma)$ . The parameters

describing an individual Gaussian peak are less certain when two or more peaks are close in energy and overlap.

# 4.2. Theoretical methods

The Amsterdam density functional theory suite (ADF 2006.01, using the standard parameters except for the options given in parentheses) was used to study the electronic structure of  $Cp_2MCl_2$  (M = V, Nb, and Ta) and Cp<sub>2</sub>VCl. The optimized geometries of Cp<sub>2</sub>MCl<sub>2</sub> and  $Cp_2VCl$  were constructed in  $C_s$  symmetry using the crystal structures as a starting point. In  $C_s$  symmetry, the Cp rings were staggered and the mirror plane is coincident with the xy plane, which bisects and each Cp ring and the Cl–V–Cl angle of Cp<sub>2</sub>VCl<sub>2</sub> or lies along the V–Cl bond of Cp<sub>2</sub>VCl. Calculations on the ground state and excited state ions were conducted in the spin unrestricted mode (unrestricted) since Cp<sub>2</sub>VCl and Cp<sub>2</sub>MCl<sub>2</sub> contain two and one unpaired electrons, respectively. A generalized gradient approximation, with the correlation of Perdew et al. [37] and exchange correction of Handy and Cohen [38] (GGA OPBE), was used. This functional has been found by us [21,39–41] and others [42–45] to give reasonable geometry parameters, spin states, and ionization and oxidation potentials. The calculations employed Slater-type orbitals for basis sets with double-zeta valence plus polarization functions for main group elements and triple-zeta valence plus polarization functions for the metal (DZP for C, H, and Cl, and TZP for V, Nb and Ta). Geometry optimizations were also carried out using integration to six significant figures (integraton = 6.0), using the zeroth-order relativistic approximation (ZORA), using tighter convergence criteria for the energy, bond distances and angles (E = 0.0005, grad = 0.001, rad = 0.0005), using the smoothing of gradients (aggressive) option, and the convergence of the selfconsistent field was tightened (converge 1e-6 1e-6).

 $\Delta$ SCF calculations of the ionized states were performed at the fixed geometry of the neutral molecule, with one electron removed from the relevant orbital. For the ligand-based ionizations both singlet and triplet ionized states for Cp<sub>2</sub>VCl<sub>2</sub> and doublet and quartet ionized states for  $Cp_2VCl$  were calculated. The  $\Delta SCF$  estimate of the ionization energy is the difference between the calculated total energy of the ionized state and that of neutral ground state molecule. As is customary for larger polyatomic molecules with low frequency vibrational modes, a semiclassical treatment of the ionization intensity is assumed that neglects zero-point vibrational energies (ZPEs). This semi-classical interpretation is reasonable since the low-frequency vibrational modes give an essentially continuum, vibrationally unresolved band in the photoelectron spectrum.

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